

Advanced Quantum Field Theory: Modern Applications in HEP, Astro & Cond-Mat
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Handout 2 (Spring 2016 term)

1. Within the theory of a free massless scalar $D = 1 + 1$, consider the operator

$$:\hat{T}_{00}:(f) \equiv \frac{1}{2} \int :(\partial_t \hat{\varphi}(x, 0))^2 + (\partial_x \hat{\varphi}(x, 0))^2 : f(x) dx,$$

where $f(x)$ is a non-negative infinitely differentiable function with a compact support (a *bump function*). Find the norm of the state $:\hat{T}_{00}:(f)|0\rangle$ and the expectation value of \hat{A} in the state

$$|\Phi_\alpha\rangle \equiv \mathcal{N}(1 + \alpha : \hat{T}_{00} : (f)) |0\rangle, \quad \alpha \in \mathbb{C},$$

where \mathcal{N} is the normalization constant, in terms of the Fourier image of f . Using explicit instances of $f(x)$, show that the density expectation value in $|\Phi_\alpha\rangle$

$$\lim_{\text{supp} f \rightarrow \{x\}} \frac{\langle \Phi_\alpha | : \hat{T}_{00} : (f) | \Phi_\alpha \rangle}{\int f(x) dx}.$$

is not positive-definite and is unbounded below.

2. Within the problem of an accelerated mirror in a $D = 1 + 1$ massless scalar theory, start with the integral expression for the regularized stress-energy tensor $\langle T_{\mu\nu}(x_-) \rangle$ given in lecture 7 and check the formulae for the renormalized stress-energy

$$\langle T_{00}^{\text{ren}} \rangle = -\langle T_{01}^{\text{ren}} \rangle = \frac{1}{12\pi} \sqrt{p'(x_-)} \left[\frac{1}{\sqrt{p'(x_-)}} \right]''.$$

3. Show that, when the effect of the external magnetic field on the photon dispersion is neglected, the photon splitting to more than 2 photons is kinematically forbidden, namely, the corresponding phase volumes vanish. Moreover, find the double-splitting momentum integral

$$V_{\text{ph}}^{(2)}(\mathbf{k}) = \int d^3 k_1 d^3 k_2 \delta^3(\mathbf{k}_1 + \mathbf{k}_2 - \mathbf{k}) \delta(|\mathbf{k}_1| + |\mathbf{k}_2| - |\mathbf{k}|).$$

4. Using the well-known expression for the Heisenberg–Euler effective action, find the term in the Lagrangian that is responsible for the photon splitting in a strong external magnetic field \vec{B} (more precisely, for the contribution of the hexagon diagram to the amplitude).
5. Within a scalar $D = 1 + 1$ theory with the Lagrangian

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \varphi)^2 - \frac{\lambda}{4} (\varphi^2 - v^2)^2, \quad m \equiv v\sqrt{\lambda},$$

find the mass counterterm $\delta\mathcal{L} = -\frac{1}{2}\delta m^2 \varphi^2$ which renormalizes the self-energy, within the leading order in λ . In the integrals, use the momentum cutoff $|p| < \Lambda \rightarrow \infty$ and explicitly expand the counterterm in divergent powers of Λ .

6. Prove the Fierz identities ($\varphi, \chi, \psi, \omega$ are classical spinor fields)

$$\bar{\varphi} \gamma_L^\mu \chi \bar{\psi} \gamma_{\mu L} \omega = \bar{\varphi} \gamma_L^\mu \omega \bar{\psi} \gamma_{\mu L} \chi, \quad \bar{\varphi} \gamma_R^\mu \chi \bar{\psi} \gamma_{\mu L} \omega = -2\bar{\varphi} P_L \omega \bar{\psi} P_L \chi,$$

where $P_{L,R} \equiv (1 \mp \gamma_5)/2$, $\gamma_{L,R}^\mu \equiv \gamma^\mu P_{L,R}$.

7. Starting with the general expression for the 1-loop effective potential for axions in the AWZ model (see lecture 11), use the momentum cutoff regularization $|p| < \Lambda \rightarrow \infty$ and find the divergent and finite parts of the potential.